Using fireball networks to track more frequent reentries: Falcon 9 upperstage orbit determination from video recordings

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ABSTRACT

On February 16, 2021, an artificial object moving slowly over the Mediterranean was recorded by the Spanish Meteor Network (SPMN). Based on astrometric measurements, we identified this event as the reentry engine burn of a SpaceX Falcon 9 launch vehicle's upper stage. To study this event in detail, we adapted the plane intersection method for near-straight meteoroid trajectories to analyze the slow and curved orbits associated with artificial objects. To corroborate our results, we approximated the orbital elements of the upper stage using four pieces of "debris" cataloged by the U.S. Government's Combined Space Operations Center. Based on these calculations, we also estimated the possible deorbit hazard zone using the MSISE90 model atmosphere. We provide guidance regarding the interference that these artificial bolides may generate in fireball studies. Additionally, because artificial bolides will likely become more frequent in the future, we point out the new role that ground-based detection networks can play in the monitoring of potentially hazardous artificial objects in near-Earth space and in determining the strewn fields of artificial space debris.

KEYWORDS

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1 Introduction

Interplanetary meteoroids generate fireballs when penetrating the Earth's atmosphere in a range of velocities between 11.2 km/s, which is the minimum velocity required for attraction by the Earth, and 73 km/s, which is the maximum velocity that a natural body gravitationally bound to our solar system can achieve [1]. These luminous phenomena are generated by large meteoroids on the centimeter or meter scale impacting the atmosphere at hypersonic velocities and ablating their components as they collide with air particles and heat up [2–4]. Some of these bodies undergo catastrophic disruption and disintegrate completely, whereas others survive atmospheric entry and are deposited on the Earth's surface. These surviving materials, if they are pristine rocks of natural origin, are called meteorites [5].

To gain a better understanding of these events, meteor detection networks have been developed worldwide to monitor the sky and obtain valuable information regarding the characteristics, fates, and origins of meteoroids [6–15]. Since 1995, the Spanish fireball and meteorite recovery network (SPMN) has been operating in Spain and is distributed throughout the peninsular and insular territory with 30 ground-based stations equipped with all-sky charge-coupled device cameras and wide-field video systems [9, 16, 17]. The SPMN records the sky on a full-time basis and automatically detects any moving luminous objects up to a magnitude of 10.

However, not all fireballs recorded by these detection

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systems have natural origins. Some luminous events are generated by artificial meteors created by human space programs and the proliferation of satellite technology. The ever-increasing number of new launches, collisions in space, and debris-shedding events generate space debris in low Earth orbits [18]. This technogenic pollution generates a two-fold problem: (i) it poses a risk for space exploration and can cause damage when debris fall back to the ground [19], and (ii) it can interfere with astronomical observations [20]. In this regard, space debris in low orbits undergo gradual decay induced by atmospheric drag until they eventually burn up or flare by reflecting sunlight in favorable geometries. These events are largely undesirably recorded by optical systems [21, 22]. The detection and analysis of such events can contribute to monitoring the deorbits of hazardous artificial objects in near-Earth space.

Although artificial object reentries are eventually analyzed through ground-based observations (e.g., the simple-return capsule Genesis [23], robotic space probe Stardust [24, 25], cargo spacecraft Jules Verne ATC [26], and asteroid explorer Hayabusa [27, 28]), no meteor detection networks have developed and systematically implemented detection and reduction algorithms specifically for artificial meteor analysis.

Given the increasing number of such events and the need to discern their origins and fates, we adapted the SPMN network reduction method, which is available in the new 3D-FIRETOC software [29], for the automatic detection and analysis of objects with slow and curved trajectories. We tested our implementation through the study of the Falcon 9 reentry and compared the results to calculations based on debris orbital elements. Finally, we estimated the possible deorbit hazard zone using the MSISE90 atmospheric model.

2 Dataset analyses

We present two analyses of the event labelled SPMN160221ART recorded by the SPMN on February 16, 2021, which crossed southern France in the direction of Libva. The fireball was captured by the north and east cameras of the Estepa station in Seville province, as well as the eastern camera of Benicàssim station in Castellón province (see Table 1). The object could be observed in the field of view for a minute and a half. which indicated a low velocity and elliptical or nearly circular orbit. As will be demonstrated in this paper, this event was generated by the upper stage and payloads of a Falcon 9 rocket that launched a batch of SpaceX Starlink satellites into Earth orbit earlier that night. This rocket stage was deliberately deorbited over the Indian Ocean 1.5 revolutions (2.5 h) after launch (just after our observations). Figure 1 presents enhanced images extracted from each video recording. When comparing the images to those captured by the International Space Station and images of Iridium flares, we concluded that the artificial fireball exhibited an apparent magnitude of -6 ± 1 at both monitoring stations.

2.1 Astrometric calibration

To reduce both recordings astrometrically, the first step is to perform astrometry on the stars recorded in the images containing the artificial object/bolide to calculate its projection onto the celestial sphere (i.e., its apparent trajectory) [29]. By identifying the stars in the visible sky, a method can be applied to consider the distortion of the lens to determine the relationships between pixels and horizontal coordinates. Various calibration methods have been proposed for all-sky cameras [30–32]. These calibrations involve solving highly nonlinear equations such that convergence is nontrivial. Therefore, we implemented the polynomial variant proposed by Ref. [33], which significantly improves the convergence of solutions. However, based on the low number of stars visible in the videos and unknown camera constants, we applied the simplex algorithm to optimize the initial values and guarantee a robust solution [34]. The proposed method follows the diagram presented in Fig. 2, assumes symmetrical lens distortion, and requires the determination of the parameters P_1 , P_2 , P_3 , a_0 , E, ϵ , x_0 , and y_0 according to the following equations:

Table 1 SPMN stations recording the SPMN160221ART event on 2/16/2021

Station	Longitude	Latitude	Alt.	Start time	End time
Estepa N	$4^{\circ}52'36'' W$	37°17′29″ N	537 m	05:53:36 UTC	05:56:30 UTC
Estepa E	$4^{\circ}52'36'' W$	37°17′29″ N	537 m	05:55:53 UTC	05:58:51 UTC
Benicàssim	$0^{\circ}02'19'' E$	40°02′03″ N	15 m	05:57:38 UTC	05:59:03 UTC





Fig. 1 Top left: classic S-shaped cloud from the Estepa North video with false color enhanced. Top right: Falcon 9 upper stage, payload swarm, and expanding fuel/gas cloud from Estepa North. Bottom left: overlaid frames and reference stars from Benicàssim. Bottom right: overlaid frames and reference stars from Estepa East.

$$r = \sqrt{(x - x_0)^2 + (y - y_0)^2} \tag{1}$$

$$u = P_1 r^2 + P_2 r + P_3 \tag{2}$$

$$b = a_0 - E + \tan^{-1} \left(\frac{y - y_0}{x - x_0} \right)$$
(3)

$$\cos(z) = \cos(u)\cos(\epsilon) - \sin(u)\sin(\epsilon)\cos(b) \qquad (4$$

$$\sin(a - E) = \sin(b)\sin(u)/\sin(z) \tag{5}$$

where x_0, y_0 is the center of projection (COP), where the system's optical axis intersects the sensor plane; r, u, and b are the radial distance, zenith-like angle mapping, and azimuth-like angle of a pixel coordinate and the COP, respectively; a_0 is the rotation of the sensor's x axis from the cardinal south; E is the rotation between the x axis and a vector defined by the true zenith projection and the COP; ϵ is the angle between the true zenith and COP; and z and a are the zenith angle and azimuth of the given pixel coordinate, respectively.

Once the cameras are calibrated, it is possible to transform pixel coordinates into horizontal coordinates and then equatorial coordinates. Based on the projections of the apparent trajectory onto the celestial sphere, the real trajectory can be reconstructed. However, because the time span of the video is relatively long, the rotational motion of the Earth is relevant and must be considered when transforming between coordinate systems.

2.2 Orbit reconstruction from ground-based observations

We originally implemented the triangulation method based on the intersection of planes proposed in Ref. [30]. However, to study this event, we had to adapt this method to computing curved paths such as satellite orbits because our implemented meteor analysis software was specifically designed to compute typical near-straight trajectories of meteoroids [35].

The mean plane containing the apparent trajectory of each station was obtained from the observation of a fireball from two or more stations. The intersection of these planes represents the atmospheric trajectory of the meteoroid. To reconstruct a curved orbit, we divided the observed trajectory into small segments such that each segment could better conform to a linear assumption (see Fig. 3). In this manner, we obtained small straight sections that formed the curved trajectory when combined. The division into small segments was performed in accordance with the total number of observed points, which



Fig. 2 Block diagram of the simplex method applied to the astrometry. R is the substitution point for reflection, E is the expansion, C is the contraction, and S is the shrinkage.



Fig. 3 Schematic diagram of the plane intersection method divided into segments to reconstruct curved trajectories considering the rotational motion of the Earth.

corresponded to an equally distributed duration because during the observation period, the velocity change was sufficiently slow (less than the uncertainty of the observed velocity). Because the plane intersection method corrects for point spread by calculating the mean plane of a path, overly small segments would be unable to compensate for deviations, whereas overly large segments would not correctly model the curved trajectory. Therefore, given the sensitivity of the measurements, by reducing the size of each segment and increasing the number of divisions, the triangulation method will begin to diverge at some point. As shown in Fig. 4, the results are robust below eight subdivisions.

By fitting the mean plane containing the observed points (the plane to which they are the least distant on average), we calculated the plane of the orbit. Points that did not lie on this plane were projected perpendicularly onto the plane to minimize error. In this manner, we derived the corresponding Cartesian coordinates of each detected point such that the orbital state vector at a given epoch could be trivially derived by subtracting two consecutive positions. We achieved the best fit with four segments, resulting in minimized mean residuals. Figure 5 presents the residuals of the fitted elliptical orbit.

The last step was to transform the state vector into

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Fig. 4 Variation in mean velocity, inclination, argument of perigee, and right ascension of the ascending node (RAAN) as a function of the number of segments used in the modified plane intersection method.



Fig. 5 Residuals of the elliptical orbit fitted by subdividing the observed path into four segments.

orbital elements. Once we had the position and velocity vectors, we calculated the specific angular momentum \bar{h} and node vector \bar{n} as follows:

$$\bar{h} = \bar{r} \times \bar{v} \tag{6}$$

$$\bar{n} = (-h_y, h_x, 0) \tag{7}$$

Then, all the orbital elements could be computed as

$$a = \frac{1}{\frac{2}{r} - \frac{v^2}{GM}}\tag{8}$$

$$\bar{e} = \frac{1}{GM} \left[\left(v^2 - \frac{GM}{r} \right) \bar{r} - (\bar{r} \cdot \bar{v}) \bar{v} \right]$$
(9)

$$i = \arccos\left(\frac{h_z}{h}\right) \tag{10}$$

$$\Omega = \arccos\left(\frac{h_y}{\sqrt{h_x^2 + h_y^2}}\right) \tag{11}$$

$$\omega = \arccos\left(\frac{-h_y e_x + h_x e_y}{e_\sqrt{h_x^2 + h_y^2}}\right) \tag{12}$$

$$v_0 = \arccos\left(\frac{\bar{e} \cdot \bar{r}}{er}\right) \tag{13}$$

where a is the semi-major axis, e is the eccentricity, i is the inclination, Ω is the longitude of the ascending node, ω is the argument of the perihelion, and v_0 is the true anomaly [36].

2.3 Proxy orbit computation from debris piece tracking

After identifying the event as potentially being related to the Starlink V1.0-L19 launch on the same night, we obtained orbital elements for the objects from the U.S. Government's Combined Space Operations Center (CSpOC) through their web portal Space-Track. For this launch, 64 objects were cataloged: 60 payloads and four pieces of debris. Table 2 lists the IDs of the four debris objects



considered				
Debris	NORAD ID	COSPAR ID		
1	47683	2021-012BR		

2021-012BQ

2021-012BP

2021-012BN

47682

47681

47680

Table 2 NORAD and COSPAR IDs of the debris pieces

considered. The pieces of debris are four containment rods that were used to keep the payloads stacked on the Falcon 9 upper stage. They were jettisoned during payload release.

Because objects that make less than two revolutions before reentry are typically not cataloged, there are no tracking-based orbital elements for the Falcon 9 upper stage, which is the prime suspect for the event we observed. However, we can derive approximate orbital elements for the upper stage by using the four pieces of debris as a proxy. The debris were jettisoned upon payload release under the effects of inertia and subjected solely to drag (unlike the payloads, which subsequently maneuvered to higher orbits using their propulsion systems). The orbit of the Falcon 9 upper stage should initially closely match those of the four debris pieces. The first elements available for the four debris pieces were captured on February 23, 2021, which was one week after launch. By using the SGP4 computational model [37, 38], we propagated the orbits back to the moment at which they separated from the upper stage and satellite stack (February 16, 2021, 04:08:24 UTC). We then took the average of the four resulting orbital element sets as the first proxy for the orbit of the Falcon 9 upper stage. For the two observation stations at Estepa and Benicàssim, this first proxy orbit resulted in sky trajectories that closely match the observations with a small time difference Δt of ~ 17 s for a given point on the sky trajectory. Small tweaks to the mean motion, eccentricity, and inclination were then made to reduce Δt to near zero, resulting in a new final proxy orbit.

We employed the SatFit 3.1 orbit fitting software written by Scott Campbell for this process^①. SatFit modifies the SGP4 orbital elements by using a least-squares fitting procedure to improve the fit of elements to astrometric observations. SatFit provides feedback regarding the resulting fit by returning information on the overall posi-

tional error, cross-track error, and Δt . The mean motion, which represents the time it takes the rocket stage to complete one revolution around the Earth, was adjusted to a value yielding a Δt of less than 1 s via fitting to four astrometry points from the early portion of the Estepa video. The inclination was adjusted to eliminate a small cross-track error. The adjustments amounted to -0.00009242 rev/day for mean motion and 0.08° for inclination. Small adjustments were also made in the RAAN (-0.3029°) , eccentricity (-0.0000026), and mean anomaly to improve the Δt fit and reduce the cross-track error (adjusting the inclination automatically means the RAAN must also be adjusted and adjusting the mean motion means the mean anomaly must also be adjusted).

2.4 Reentry trajectory prediction

The triangulation of the bright phase in the images from Estepa and Benicàssim facilitated the construction of a state vector and from this vector, a second set of orbital elements was derived (Table 3). The state vector was transformed into orbital elements using the RV2TLE software written by Scott Campbell, yielding a set of SGP4compatible orbital elements in the three-line-ephemerid (TLE) format².

The proxy orbit of the Falcon 9 upper stage is consistent with the reduction performed on the ground-based video data, as can be observed in Fig. 6, particularly in the orbital plane, but also in the orbital altitude.



Fig. 6 3D scale representation of the Falcon 9 trajectory orbit. The proxy orbit is white (yellow when illuminated by sunlight), the observed path used for reduction is purple, the Estepa station is green, and the Benicàssim station is orange.

Our triangulation calculations yielded an altitude of approximately 270 ± 0.6 km and velocity of 7.5 ± 0.3 km/s

 $\mathbf{2}$

3

4

① The source code is available at http://sat.belastro.net/satelliteor bitdetermination.com/.

⁽²⁾ This format is commonly used for satellite orbits (for a description of the TLE format, see http://www.satobs.org/element.html).

Table 3	TLE computed t	from the	triangulation	of SPMN data
	1			

Falcon 9 R/B (from State Vector)					
1 99999U 21012BS 21047.24849537 0.00000000 00000-0 00000+0	0 09				
2 99999 52.4033 104.0553 0228149 225.2325 266.6565 16.052387	01 08				

at 05:57:58 UTC. It should be noted that our observations likely captured the rocket stage just after it performed a deorbit burn. Therefore, small discrepancies are expected between our first orbital element set, the proxy orbit (which is the pre-burn orbit), and real observed trajectory. However, at this stage of the launch, the payloads should still be very close to the pre-burn proxy orbit for the upper stage. The second orbital element set derived from the state vector represents the post-burn reentry trajectory. This orbit has a semi-major axis of 6638 km with a nominal apogee of 411 km and perigee of 108 km. These are values with respect to the Earth's equatorial radius. The real altitude depends on the location of the perigee (for this particular orbital revolution, the perigee is near 118 km above the geoid). The orbital inclination is 52.4° , which is a difference of a few tenths of a degree (hundred arcseconds) relative to the pre-burn proxy orbit discussed above. With an eccentricity of 0.0228, this orbit is much more eccentric than the pre-burn proxy, as expected for a post-burn reentry trajectory.

The perigee on this revolution is reached at approximately 6:19 UTC and is located near 30.1° S, 68.9° E. This is within the western portion of the deorbit hazard zone defined by Navigational Warning HYDROPAC $463/21^{\circ}$. The map in Fig. 7 compares the orbit derived from the state vector (solid white line) to the pre-burn proxy orbit (thin dashed line). The yellow cross represents the perigee. Time is in UTC.

Compared to the pre-deorbit-burn proxy orbit based on the orbits of the four retaining rods, the post-deorbitburn orbit from the triangulation-based state vector has a difference of -0.732 deg in inclination, -11.28 km in the semi-major axis, and 0.020423 in eccentricity. The eccentricity of the post-reentry burn orbit must be larger than the eccentricity of the pre-deorbit-burn proxy orbit by definition. Its semi-major axis must also be smaller by definition.

The three-dimensional (3D) positional difference between the two orbits is minimized (10 km) at approximately 5:57:48 UTC. The cross-plane difference is mini-



Fig. 7 Reentry trajectory estimation. The proxy orbit is the dashed white line and the triangulated orbit is the solid white line. The yellow cross marks perigee of this orbit. The area indicated in red is the deorbit hazard area identified by the Navigational Warning HYDROPAC 463/21.

mized (5 km) 45 s earlier at approximately 5:56:58 UTC. At the state vector epoch (05:57:50 UTC), the absolute positional difference between the pre-deorbit-burn proxy orbit and state vector is approximately 10.2 km, 1.8 km of which is in altitude and 8.4 km is in the cross-plane (horizontal) direction. Compared to the inherent positional accuracies of SGP4 (1 km at epoch and increasingly more before and after that), these differences are small and reasonable differences indicating a good fit between the proxy orbit and orbit derived from triangulation [39, 40].

The moment of the smallest positional difference (5:57:48 UTC) is very close to the state vector epoch, which is close to the moment at which the trail on the images noticeably brightened. This may indicate that this moment captures the start of the actual deorbit burn. Because the deorbit burn has a duration, this could indicate that the state vector we obtained still underestimated the eccentricity and overestimated the perigee altitude of the final reentry orbit.

The nominal perigee altitude derived from the state vector appears to be slightly too high. A reentry model was simulated using the NASA General Mission Analysis Tool (GMAT) R2020a software $[41]^2$ with the MSISE90

 $[\]textcircled{O}$ GMAT is downloadable at https://sourceforge.net/projects/gma t/.



① https://msi.nga.mil/NavWarnings.

model atmosphere, nominal orbit from Table 3, space weather at that time, and dry mass and drag surface values for a Falcon 9 upper stage (4500 kg and 58.5 m² maximum drag surface). The model indicated that the nominal orbit derived from the state vector should see the rocket stage survive perigee and continue for a few additional revolutions.

However, an orbit that does result in deorbit over the designated area is possible within the error margins of the speed vector. A reduction in the speed vector of only 0.015 km/s brings the perigee altitude of the orbit to below 80 km and modeling in GMAT then results in a deorbit inside the designated zone from the Navigational Warning HYDROPAC 463/21. Modeling was performed twice: once for a maximum drag surface of 58.5 m^2 and once for a minimum drag surface of 10.5 m^2 . The rocket-stage dry mass used in both cases was 4500 kg. The resulting modeled impact points were near 38.0 S. 79.2 E (maximum drag surface scenario), and 49.3 S. 104.3 E (minimum drag surface scenario), which are both within the area indicated by the Navigational Warning HYDROPAC 463/21. Furthermore, as indicated earlier, our state vector may actually represent the start of the deorbit burn, so it may not capture the full effect of the burn on the final reentry orbit, which could lead to an overestimation of the perigee altitude.

Based on these observations, practically no acceleration can be inferred. In the observed path, the velocity change ranges from 7.4499 to 7.4644 km/s (from the proxy orbit). However, from the ground stations, it is difficult to obtain high-accuracy velocity measurements (below 0.1 km/s). We observed a velocity of 7.5 ± 0.3 km/s, which does not allow us to appreciate the speed changes. Considering the fact that the recording contains the pre-ignition fuel cloud, subsequent engine burn, and consistency between the debris pieces, orbits, the Falcon 9 Starlink V1.0-L19 launch and our observations are evidence that the SPMN160221ART event was a controlled deorbiting maneuver in a quasi-circular orbit.

3 Discussion

It seems clear that the increasing use of the near-Earth environment for commercial purposes will make reentries more frequent events (even with an increasing practice of deliberate deorbiting over the southern Pacific Ocean at the end of service life). In fact, another reentry was widely observed over the USA on March 26, 2021. In that case, the object was a Falcon 9 upper stage from the March 4, 2021 Starlink launch. The Falcon 9 upper stage failed to deorbit for unknown reasons and came down uncontrolled over Oregon and Washington in the northwest of the USA. This attracted significant public attention and several casual eyewitnesses filmed the event using their mobile phones. These recent events exemplify why we should be able to recognize and explain the real nature of such appearances. For the aforementioned reasons, it is relevant to develop a common methodology and software solution.

Regarding the event on February 16, 2021, which is the subject of this study, we should discuss the sequence of events of typical controlled rocket stage reentries to understand our observations. A deorbit burn has at least two phases: an actual engine burn and propellant blowout (fuel vent) at the end of the burn to avoid explosive disintegration of the rocket stage. Fuel vents (propellant blow-outs) often generate a circular cloud, sometimes (particularly with Falcon 9 stages, where it has been reported several times) with a spiral shape if the stage is rotating (spin stabilization). The fuel cloud initially moves at the same speed as the rocket stage from which it originates (the cloud co-orbits with the rocket stage). Over time, it expands, and differential drag and small ΔV differences separate it from the rocket stage. However, just after blow-out, it will stay with the rocket stage for a short duration as it begins radially expanding away from the rocket stage [42]. This is what we observed in our records of the February 16 event. In the early part of the Estepa North camera record, starting just after 5:53:45 UTC, a diffuse circular cloud can be observed surrounding one of two faint objects (see Fig. 1). This is approximately 4 min before the start of the sudden bright phase. One of the two faint objects is likely the clump of released payloads. The other, which is centered in the diffuse cloud, is likely the Falcon 9 upper stage. This suggests that a burn or propellant blow-out (tank depressurization) occurred just before the start of camera recording. The beginning of visibility in the Estepa North camera record corresponds to the time at which rocket stage and generated fuel cloud passed from the Earth's shadow into sunlight at approximately 5:53:45 UTC.

These findings make it very unlikely that the sudden bright phase appearing during the pass imaged by both Benicàssim and Estepa was caused by ablation. Therefore, we prefer the interpretation that the bright phase represents the actual deorbit engine burn. Alternatively, it could be generated by the payloads flaring up when the Sun-payload-observer angle and the angles of the payload surfaces are favorable. The payloads may be very bright occasionally and in this phase of the mission, they are still close to the rocket stage (as shown at the start of visibility in the Estepa video). The pre-burn proxy orbit and the orbit derived from the state vector resulting from the observations converge to within 10 km at 5:57:48 UTC. This is very close to the start of the sudden bright phase. The fuel cloud captured a few minutes earlier could be attributed to an earlier maneuver (e.g., a payload avoidance maneuver).

4 Conclusions

It is becoming increasingly common to observe luminous objects moving through the sky. Although most fireballs are natural and have a meteoric origin, a growing number of human-made objects in orbit around the Earth is leading to countless sightings and detection records of impressive light phenomena in the sky. In this work, we demonstrated how a traditional fireball analysis technique can be adapted to compute the curved trajectories of objects experiencing deorbit and reentry, as well as how to handle long recording durations.

- We developed an extension for the 3D-FIRETOC Python software by modifying the plane intersection method for meteor triangulation to be able to analyze slow objects with curved orbits captured over long periods of time.
- We computed a Falcon 9 upper stage trajectory from video recordings obtained by two SPMN network stations on February 16, 2021. The results were successfully compared to an orbit estimate from the orbital parameters published by CSpOC for four debris pieces associated with the Starlink V1.0-L19 launch.
- Our data confirmed that the recordings were obtained during or just after the deorbit burn of the rocket stage, meaning this type of pre-ablation phase in a reentry has the potential to produce eyewitness sightings causing varying degrees of alarm because such sightings are typically unexpected.
- By using the 3D-FIRETOC pipeline, the final orbit and deorbit trajectory of the Falcon 9 stage were successfully modeled. The results indicated a reentry over

the area in the Indian Ocean that was designated for this entry, which is an additional added value of reentry tracking from the ground that increases our capacity to recover space junk arriving on the ground.

- It seems clear that the advancement of space technology will lead to an increase in satellite sightings and artificial bolides produced by the reentry of space debris and rocket stages. This may interfere with astronomy in general and fireball studies in particular.
- Based on the large amount of data recorded and the proliferation of new detection stations around the world, fireball analysis processes are being automated. Therefore, the development of false-positive avoidance techniques to rule out unnatural events of no scientific interest is urgent.
- This type of event will become increasingly common and our results exemplify how ground-based fireball station networks could gradually play an important role in monitoring the deorbit of hazardous artificial objects in near-Earth space.

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In 2003, he received a USA–Spanish grant that allowed him to continue his career in a postdoctoral position at the Institute of Geophysics & Planetary Physics of the University of California Los Angeles (UCLA) and the NASA Astrobiology Center at UCLA under the supervision of Professor John Wasson and Doctor Alan Rubin. After almost three years of work on the transport of water and volatiles in primitive meteorites (carbonaceous chondrites), he returned to Spain in 2006 with a Juan de la Cierva grant to join the Institute of Space Sciences (ICE, CSIC-IEEC) in Barcelona, Catalonia. In 2009, he gained his position as a tenured scientist of the Upper Research Council (CSIC) at the same research institute. Since 2010, Dr. Trigo-Rodriguez has been the leader of the Mete-

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Security Center of the Royal Dutch Air Force. He is still affiliated as a guest researcher at the Naturalis Biodiversity Center. He received the Van Es Prize for Dutch Archeology in 1998 and the Doctor J. van der Bilt Prize of the Royal Dutch Association for Meteorology and Astronomy (KNVWS) in 2012. In 2008, the IAU named the asteroid (183294) Langbroek in his honor. He is active as a popular science educator, including appearances in news media and on Dutch radio and television on topics such as meteorites, fireballs, and satellites. E-mail: Macro@langbreek.org.



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